

NEPOLARIZIRANI PRESEK

$$e^-\bar{\mu} \rightarrow e^-\bar{\mu}$$

$s_1 \quad s_2 \quad s_3 \quad s_4$

s_1, s_2 : POLARIZIRANI e, μ
 s_3, s_4 : NEVZIM POLARIZACIJO
 (SILKE SPINA)

NEPOLARIZIRAN PRESEK: SIPALNI PRESEK ZA NEPOLARIZIRANE VPADNE DELCE IN ZA PRIMER, KO NE VZEMIM SPINOV V KONONEM STANJU.

$$\mathcal{Z} \propto |\mathcal{M}|^2$$

POLARIZIRAN PRESEK

$$\mathcal{Z} \propto |\mathcal{M}_{s_1 s_2 s_3 s_4}|^2$$

NEPOLARIZIRAN PRESEK

$$\mathcal{Z} \propto \frac{1}{(2s_a+1)(2s_b+1)} \sum_{s_1 s_2 s_3 s_4} |\mathcal{M}_{s_1 s_2 s_3 s_4}|^2$$

s_a, s_b ZA
 NAŠ PRIMER:
 $s_a = \frac{1}{2}, s_b = \frac{1}{2}$

$$\frac{1}{(2s_a+1)(2s_b+1)} = \frac{1}{4}$$

ZAVAJ

$$\sum_s |\mathcal{M}_{s_1 s_2 s_3 s_4}|^2 \cdot \text{IN NE } \left| \sum_s \mathcal{M}_{s_1 s_2 s_3 s_4} \right|^2$$

ZATO, ker LAHKO SPINE V ZADETKU
 IN NA KONCU V PRINCIPU IZMERIMO. SESTEVANJE
 AMPLITUD PRIDE V POSEN, KADAR PROCESI
NISO LOGLINI.

$$u^{(k)}(\ell) g^\alpha$$

$$|\mathcal{M}|^2 \propto [\bar{u}(\ell') g^\alpha u(\ell)] [\bar{u}(\ell') g^\beta u(\ell)]^* = [\bar{u}(\ell') g^\alpha u(\ell)] [\bar{u}(\ell') g^\beta u(\ell)]^*$$

$$= [\bar{u}(\ell') g^\alpha u(\ell)] [u(\ell)^+ g^{\beta+\alpha} u(\ell)] =$$

$g^{\beta+\alpha} = g^\alpha g^\beta$

$$= [\bar{u}(\ell') g^\alpha u(\ell)] [u(\ell)^+ g^\beta u(\ell)] = [\bar{u}(\ell') g^\alpha u(\ell)] [\bar{u}(\ell) g^\beta u(\ell)]$$

$$\sum_{s_1 s_2} [\bar{u}^{(s)}(\ell') g^\alpha u^{(s)}(\ell)] [\bar{u}^{(s)}(\ell) g^\beta u^{(s)}(\ell)] =$$

$$= \sum_{S,S'} [\bar{u}_{\alpha}^{(S)}(z) \gamma^{\nu c} u_{\beta}^{(S)}(z)] [\bar{u}_{\delta}^{(S)}(z) \gamma^{\nu c} u_{\epsilon}^{(S)}(z)] =$$

$$= \sum_{S,S'} \underbrace{u_{\epsilon}^{(S)}(z) u_{\alpha}^{(S)}(z)}_{[\not{k} + m_e]_{\epsilon\alpha}} \gamma^{\nu c} \underbrace{\bar{u}_{\beta}^{(S)}(z) \bar{u}_{\delta}^{(S)}(z)}_{[\not{k} + m_e]_{\beta\delta}} \gamma^{\nu c}$$

$$= [\not{k} + m_e]_{\epsilon\alpha} \gamma^{\nu c} [\not{k} + m_e]_{\beta\delta} \gamma^{\nu c} =$$

$$= [\not{k} + m_e] \gamma^{\nu c} [\not{k} + m_e] \gamma^{\nu c} = \text{Tr}[(\not{k} + m_e) \gamma^{\nu c} (\not{k} + m_e) \gamma^{\nu c}]$$

ENAKI RAZMISLEK ZA MORE

$$\overline{|M|^2} = \frac{1}{2} \cdot \frac{1}{2} \frac{e^2 e^2}{2^4} \cdot \text{Tr}[(\not{k} + m_e) \gamma^{\nu c} (\not{k} + m_e) \gamma^{\nu c}] \text{Tr}[(\not{k} + m_e) \gamma^{\nu c} (\not{k} + m_e) \gamma^{\nu c}]$$

TEOREMI, O SLEDEN

EDEN OD NJIH

$$\text{Tr}[(\not{k} + m_e) \gamma^{\nu c} (\not{k} + m_e) \gamma^{\nu c}] = 4 \left[\not{k}^{\nu c} \not{k}^2 + \not{k}^{\nu c} \not{k}^2 - (\not{k} \cdot \not{k}) \right]$$

$$= 4 \left[\not{k}^{\nu c} \not{k}^2 + \not{k}^{\nu c} \not{k}^2 - (\not{k} \cdot \not{k} - m_e^2) \right]$$

$$\overline{|M|^2} = 4 \frac{e^4}{2^4} [(\not{k} \not{p}') (\not{k} \not{p}) + (\not{k} \not{p}) (\not{k} \not{p}') - m_e^2 \not{p}' \not{p} - m_\mu^2 \not{k} \not{k} + 2 m_\mu^2 m_e^2]$$

- LOKALNE INVARIANTNE

(naj scalarni produkti)

SIPATNI PRESNI

$$\frac{d\omega}{d\Omega_i} = \frac{\partial W_{fi}/\partial\Omega_i}{\rho_i v_i}$$

$$\frac{dW_{fi}}{d\Omega_i} = \frac{2\pi}{\hbar} |T_{fi}|^2 \frac{d\rho_f}{d\Omega_i}$$

$$d^3N = V \frac{d^3p}{(2\pi\hbar)^3} = \frac{d^3p}{\rho (2\pi\hbar)^3}$$

$$\rho = \frac{1}{V}$$

NRERATIVISTIČNI
PRIMER

RELATIVISTIČNI DELCI

$$d^3N = \frac{V}{2E} \frac{d^3p}{(2\pi\hbar)^3}$$

$$d^3N = \frac{V}{2E} \frac{d^3p}{(2\pi\hbar)^3}$$

$$\rho_e \propto \frac{d^3p}{E}$$

LORENTZOVIA TRANSF.

$$dp_x' = \gamma(p_x - \beta E) \quad dp_y' = \gamma p_y \quad dp_z' = \gamma p_z$$

$$dE' = \gamma(E - \beta p_x)$$

v smjeri x

$$\frac{d^3p'}{E'} = \frac{\gamma(dp_x - \beta dE) dp_y dp_z}{\gamma(E \beta p_x)} = \frac{dp_x (1 - \beta \frac{dE}{dp_x}) dp_y dp_z}{E - \beta p_x}$$

$$= \frac{dp_x (1 - \beta \frac{p_x}{E}) dp_y dp_z}{E (1 - \beta \frac{p_x}{E})}$$

$$= \frac{d^3p}{E}$$

LORENTZOVIA INVARIJANTA

$$E^2 = p_x^2 + p_y^2 + p_z^2 + m^2$$

$$E dE = p_x dp_x$$

$$\rho_i N_i = \frac{2E_a}{V} N_a \frac{2E_b}{V} N_b$$

UPADNI DEZEC TRANS

PRIMER, KO DEZEC b
MIRUJE, O PA SE NA
NJEM SIPLJE

OE SE OBASUA GIBLJETA

$$\rho_i N_i = F = \frac{2E_a}{V} \frac{2E_b}{V} |\vec{N}_a - \vec{N}_b|$$

$$\frac{v}{c} = \beta = \frac{\gamma m v \cdot c}{\gamma m c c} = \frac{cp}{E}$$

$$\vec{N}_a - \vec{N}_b = |\vec{p}_a E_b - \vec{p}_b E_a|$$

$$E_a E_b$$

$$|\vec{p}_a E_b - \vec{p}_b E_a| = \sqrt{(p_a p_b)^2 - m_a^2 m_b^2}$$

\Rightarrow TUDI TO JE RELATIVISTIČKA INVARIJANTA

DIFERENCIJALNI PRESEK

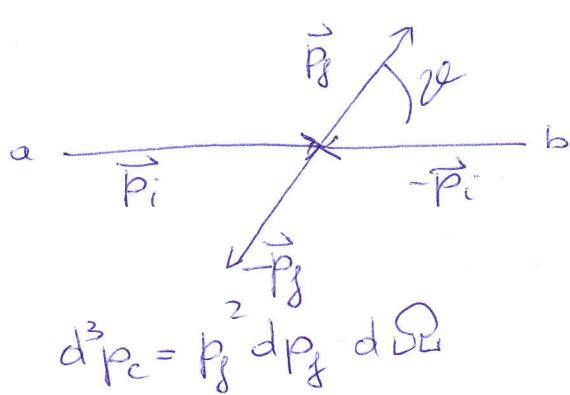
$$d\sigma = \frac{|M|^2}{F} dQ$$

ab \rightarrow cd

$$F = 4 \sqrt{(p_a p_b)^2 - m_a^2 m_b^2}$$

$$dQ = (2\pi)^4 \delta^4(p_1 + l_1 - p_2 - l_2) \cdot \frac{d^3p_c}{(2\pi)^3 2E_c} \cdot \frac{d^3p_d}{(2\pi)^3 2E_d}$$

$$-iM = \left[f_{ca}^\mu \right] \left(\frac{g_{\mu\nu}}{2} \right) \left[j_{db}^\nu \right]$$



$$\vec{P}_c = \vec{p}_c = -\vec{p}_d$$

$$\vec{P}_d = \vec{p}_d = -\vec{p}_c$$

$$d^3 p_c = p_f^2 d p_f d\Omega$$

$$d\Omega = 2\pi \sin\theta d\varphi$$

$$\int \frac{d^3 p_d}{2E_d} \underbrace{\delta(p_c + p_d - p_e - p_b)}_{\delta(E_c + E_d - E_a - E_b) \delta^3(\vec{p}_c + \vec{p}_d - \vec{p}_e - \vec{p}_b)} = \frac{1}{2E_d} \delta(E_c + E_d - E_a - E_b)$$

ENERGIA v TEZÍSNEM SISTEMU $E = E_a + E_b$

$$dQ = \frac{1}{4\pi^2} \frac{p_f^2 d p_f d\Omega}{4E_c E_d} \delta(E_c + E_d - E)$$

$$E = E_c + E_d = \sqrt{p_f^2 + m_c^2} + \sqrt{p_f^2 + m_d^2}$$

$$\frac{dE}{dp_f} = \frac{p_f}{E_c} + \frac{p_f}{E_d}$$

$$dp_f = dE \left(\frac{p_f}{E_c} + \frac{p_f}{E_d} \right)^{-1} = dE \frac{E_c E_d}{p_f (E_c + E_d)}$$

$$dQ = \frac{1}{4\pi^2} \frac{p_f dE}{4(E_c + E_d)} \delta d\Omega \delta(E_c + E_d - E)$$

$$\int dE \rightarrow dQ = \frac{1}{4\pi^2} \cdot \frac{p_f}{4E} d\Omega$$

$$F = 4p_i E$$

$$\boxed{\frac{d\Omega}{d\Omega} = \frac{|M|^2 p_f}{64\pi^2 p_i E^2}}$$

v ULTRARELATIVISTIČNÍ LIMITU $m_x \ll p_x$ $p_i = p_f$

$$\boxed{\frac{d\Omega}{d\Omega} = \frac{|M|^2}{64\pi^2 E^2}}$$

UPORABIMO TA REZULTAT ZA $e^- \mu^- \rightarrow e^- \mu^-$,
NEPOLARIZIRANO SIRANJE. V ULTRAREZAT. LIMITI

$$\vec{k}_2 = \left(\frac{E}{2}, \vec{p}_i \right) \quad \vec{k}'_2 = \left(\frac{E}{2}, \vec{p}'_i \right) \quad \vec{p} = \left(\frac{E}{2}, -\vec{p}_i \right) \quad \vec{p}' = \left(\frac{E}{2}, \vec{p}'_i \right)$$

$$k'_2 p = \frac{E^2}{4} + \vec{p}'_i \cdot \vec{p}_i = \frac{E^2}{4}(1 + \cos\vartheta)$$

$$k_2 p = \frac{E^2}{4} + \frac{E^2}{4} = \frac{E^2}{2}$$

$$\begin{aligned} Q^2 &= (\vec{k}'_2 - \vec{k}_2)^2 = (0, \vec{p}'_i - \vec{p}_i)^2 \\ &= -(\vec{p}'_i - \vec{p}_i)^2 = -\frac{E^2}{4}((\cos\vartheta - 1)^2 + \sin^2\vartheta) = \\ &= -\frac{E^2}{2}(1 - \cos\vartheta) \end{aligned}$$

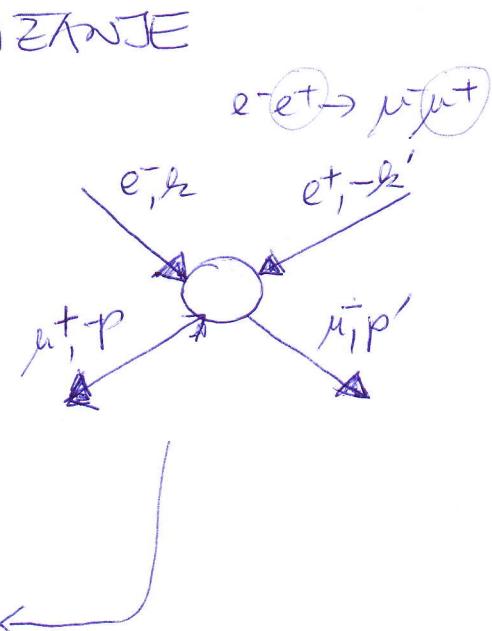
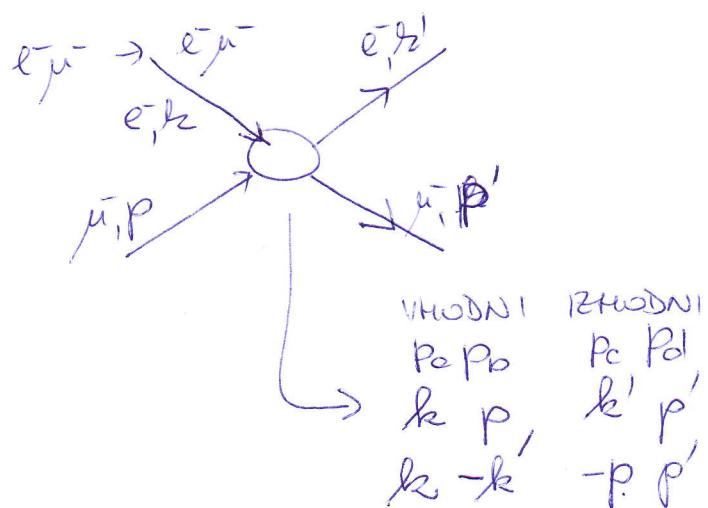
$$\begin{aligned} \vec{p}_i &= (p_i, 0, 0) = (E/2, 0, 0) \\ \vec{p}'_i &= \left(\frac{E \cos\vartheta}{2}, \frac{E \sin\vartheta}{2}, 0 \right) \end{aligned}$$

~~$\vartheta \uparrow$~~

$$\begin{aligned} |\mathcal{M}|^2 &= 4 \frac{e^4}{Q^4} [(k'_2 p)(k_2 p) + (k'_2 p)(k_2 p')] \\ &= 4 \frac{e^4}{Q^4} \left[\frac{E^4}{4} + \frac{E^4}{16} (1 + \cos\vartheta)^2 \right] \\ &= \frac{1}{4} \frac{e^4}{Q^4} [4 + (1 + \cos\vartheta)^2] E^4 \\ &= \frac{e^4}{Q^4} \frac{[4 + (1 + \cos\vartheta)^2]}{(1 - \cos\vartheta)^2} \end{aligned}$$

$$\frac{d\mathcal{M}}{dQ^2} = \frac{|\mathcal{M}|^2}{64\pi^2 E^2} = \frac{e^4 (4 + (1 + \cos\vartheta)^2)}{32\pi^2 E^2 (1 - \cos\vartheta)^2}$$

PREVERITI Doma



$$\overline{|\mathcal{M}|^2} = \frac{8e^4}{g^4} \left[(pp')(-k'k) + (k'p)(k'p') + m_e^2 p'k' + m_\mu p'k + 2m_e^2 m_\mu \right]$$

V ULTRA RELATIVISTICKÉ LIMITT

$$k = \left(\frac{E}{2}, \vec{p}_i \right), k' = \left(\frac{E}{2}, \vec{p}' \right), p = \left(\frac{E}{2}, -\vec{p}_i \right), p' = \left(\frac{E}{2}, -\vec{p}' \right)$$

$$k^2 = (k' - k)^2 \rightarrow (p + k)^2 = E^2$$

PO KREZANJU

$$\frac{d\mathcal{L}}{dk} = \frac{\overline{|\mathcal{M}|^2}}{64\pi^2 E^2} = \frac{8e^4 \left[(1 - \cos^2 \vartheta)^2 + (1 + \cos^2 \vartheta)^2 \right] \frac{E^4}{44}}{64\pi^2 E^2 E^4} =$$

$$\frac{d\mathcal{L}}{dk} = \frac{e^4}{4 \cdot 32\pi^2 E^2} \left(2 + 2 \cos^2 \vartheta \right) = \frac{e^2}{4 \cdot 16\pi^2 E^2} (1 + \cos^2 \vartheta)$$

$$\mathcal{L} = \int \frac{d\mathcal{L}}{dk} dk = \frac{e^2}{4 \cdot 16\pi^2 E^2} \int (1 + \cos^2 \vartheta) \cdot 2\pi \cdot d(\cos \vartheta) =$$

$$= \frac{e^2}{4 \cdot 8\pi E^2} \left. \left(\cos \vartheta + \frac{1}{3} \cos^3 \vartheta \right) \right|_{-1}^{+1} = \frac{e^2}{8\pi E^2} \left(2 + \frac{2}{3} \right) =$$

$$\mathcal{L} = \frac{e^2}{4 \cdot 3\pi E^2} = \frac{e^4}{12\pi E^2}$$

$$\mathcal{L} \propto \frac{1}{E^2}$$